

Estimation of Interaction Forces between Two Magnetic Bolus-like Microrobots

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Abstract—This paper analyses the interaction forces between two magnetic boluses for future drug targeting applications. To transport the drugs, it is necessary to convey several therapeutic magnetic boluses using magnetic gradients. The main difficulty is to control a group of different therapeutic boluses at desired states, despite the presence of interaction forces between boluses. To overcome this issue and designing robust control strategies, it is important to fully understand these interactions forces. Based on a dipole-dipole interaction model and dynamic modeling of two magnetic boluses, the magnetic and non magnetic forces are expressed. Finally, an experimental investigation is carried out in a tank under the presence of the magnetic field in order to assess the prevalence between the magnetic and the non-magnetic interaction forces.

I. INTRODUCTION

Magnetically actuated microrobots have been proposed for numerous applications, as their small scales enable the access to complex environments [1], [2], [3]. Especially, microrobots are applied in minimally-invasive surgery (MIS) procedure, including: targeted drug delivery, brachytherapy, hyperthermia, removing material by mechanical means or acting as simple static structures [4], [5], [6], [7]. These microrobots are commonly referred as therapeutic micro carriers (TMMC) [8], [9]. To embed the therapeutic agent, the TMMC could be either magnetic helical medical microrobots [10], [11], magnetic microbeads [12], [13], or micro/nano-particles suspended in a carrier fluid (ie. a ferrofluid) coated with organic polymer to prevent agglomeration and improve surface functionality [9], [14], [15]. Then, an external magnetic field is used to steer the TMMCs along a pre-planned path to the targeted location [8], [13]. This targeting approach is also able to steer superparamagnetic iron oxide (SPIO) particles using an improved magnetic resonance imaging (MRI) system [8], [9], [16]. To convey the desired amount of drugs, multiple boluses have to be administered, and controlled [17]. Previous studies have considered the magnetic control of a group of millimeter-sized beads immersed in fluid and driven thanks to a MRI scanner [18]. In [19] the authors have investigated the control of geometrically dissimilar Mag- μ Bots and a group of identically-fabricated microrobots. The authors proved through simulation results the stability of two millimeter-sized beads at a desired positions. However, the understanding of the complete dynamics of several microrobots remains challenging [18]. In particular, the interaction forces acting between

multiple microrobots are not fully addressed. This knowledge could greatly improve the design of robust to control law that takes into account this disturbance on the system. This study proposes to investigate these interactions forces between two microrobots. More precisely, superparamagnetic iron oxide (SPIO) particles suspended in non-magnetizable medium, termed hereafter magnetic bolus, is considered. Hence, based on a dipole-dipole interaction model, the dynamic modeling of magnetic in fluidic environment is carried out. Specifically, the magnetic and non-magnetic interaction forces between two boluses in the presence of an external magnetic field \mathbf{b}_0 is expressed. This paper is divided in four sections. Section II details the mathematical modeling of dipole-dipole interactions in order to model the magnetic interaction forces between two boluses. Then, in Section III, experiments are conducted to estimate the magnetic and non-magnetic interaction forces. Conclusion and discussions on open issues are summarized in Section IV.

II. MATERIALS AND METHODS

A. Soft Magnetic bolus Modeling

In this study, a colloidal suspension of superparamagnetic iron oxide (SPIO) particles is used as magnetic bolus. Hence, the considered magnetic microrobot is a ferrofluid droplet immersed in non-magnetizable medium. Each SPIO particles of the magnetic microrobot carry a magnetic moments. Without external magnetic field ($\mathbf{b} = 0$) their dipole directions are randomly spread, as illustrated in Figure 1(a). In such situation, the ferrofluid droplet adopts commonly a spherical shape to minimize the surface energy. In presence of an applied magnetic field density \mathbf{b} , the overall SPIO particles are polarized, and their magnetic moment are mainly aligned with \mathbf{b} . Classically, the magnetic bolus takes the shape of an ellipsoid, as depicted in Figure 1(b)-(c).

Commonly, for hard magnetic materials the magnetization \mathbf{m} is independent of the magnetic field \mathbf{b} , and could be considered easily saturated in many cases. In contrast, for soft-magnetic materials, as with using SPIO particles \mathbf{m} is strongly related to the field \mathbf{b} . At low magnetic fields, such that $|\mathbf{m}| < m_{sat}$ (with m_{sat} the saturation magnetization of the material), the magnetization of SPIO particles exhibits typically the following linear behavior [20]:

$$\mathbf{m} = \frac{\chi_a}{\mu_0(1 + \chi)} \mathbf{b} \quad (1)$$

where $\mu = \mu_0(1 + \chi)$ is the permeability of the bolus, and $\chi_a \in \mathfrak{R}^3$ is the apparent susceptibility tensor that is related to the bolus shape. Here, the considered magnetic bolus varies

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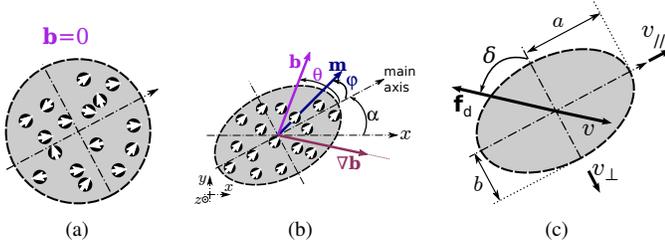


Figure 1. Soft-magnetic microrobot: (a) without external magnetic field; and (b) with an external magnetic field \mathbf{b}_0 and gradient $\nabla \mathbf{b}$ leading to an ellipsoidal ferrofluid arrangement. The black circles depict the SPIO particles with their magnetic moments \mathbf{m} shown with the white arrows. (c) Schematic model of a prolate ellipsoid of revolution ($a \geq b = c$) with a the major semiaxis and b the minor semiaxis..

from spherical geometry to axisymmetric prolate ellipsoid (see Fig. 1(c)). For spheroid the susceptibility tensor is expressed as follows [20]:

$$\chi_a = \text{diag} \left(\frac{\chi}{1 + n_a \chi}, \frac{\chi}{1 + n_b \chi}, \frac{\chi}{1 + n_b \chi} \right) \quad (2)$$

where n_a and n_b are demagnetizing factors along the major and minor axis, given by [20], [21]:

$$1 = n_a + 2n_b \quad (3)$$

$$n_a = \frac{1 - \varepsilon^2}{2\varepsilon^3} \left(\log \left(\frac{1 + \varepsilon}{1 - \varepsilon} \right) - 2\varepsilon \right) \quad (4)$$

with $\varepsilon = \sqrt{1 - \Lambda^{-2}}$ the eccentricity, and $\Lambda = a/b$ the aspect ratio of the prolate ellipsoidal magnetic bolus (see Fig. 1(c)). In the case of spherical device the demagnetizing factors simplify to $n_a = n_b = 1/3$. From the above equations, the magnetization $\mathbf{m} = (m_x, m_y, m_z)^T$ is then related to the magnetic field $\mathbf{b} = (b_x, b_y, b_z)^T$ as follows:

$$\begin{pmatrix} m_x \\ m_y \\ m_z \end{pmatrix} = \frac{V_m \chi}{\mu_0 (1 + \chi)} \begin{pmatrix} \frac{b_x}{1 + n_a \chi} \\ \frac{b_y}{1 + n_b \chi} \\ \frac{b_z}{1 + n_b \chi} \end{pmatrix} \quad (5)$$

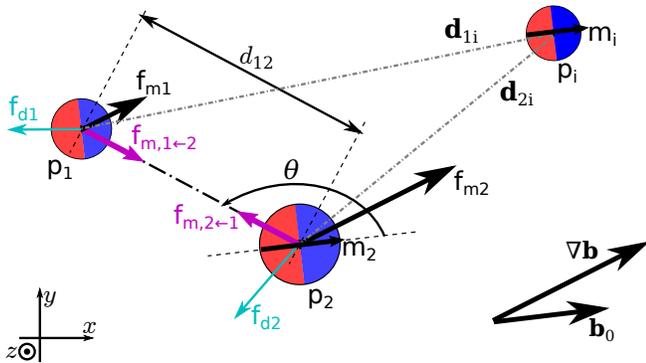


Figure 2. Group of magnetic microbeads in a microfluidic environment.

Multiple magnetic boluses have to be serially administered and controlled. A key issue in the success of the operation

is then to address the interactions forces involved in their dynamic behavior. Especially, within a magnetic field each boluses induce a dipole-dipole interaction force. It is important to study this dipole-dipole interaction force.

B. Dipole-dipole Interaction Forces

Commonly, when several magnetic particles are in a closed range, a magnetic interaction occurs. Specifically, a magnetic bolus i located at $\mathbf{p}_i = (x_i, y_i, z_i)^T$ with its dipole moment ${}^i \mathbf{m}$ induces a magnetic field in location \mathbf{p} given by [22], [23]:

$$\mathbf{b}_{\mathbf{p}_i}(\mathbf{p}) = \frac{\mu_0}{4\pi} \left(3 ({}^i \mathbf{m} \cdot \mathbf{d}_0) \mathbf{d}_0 - {}^i \mathbf{m} \right) \quad (6)$$

where $\mu_0 = 4\pi \times 10^{-7}$ T m/A is the vacuum permeability; $\mathbf{d}_0 = \frac{(\mathbf{p} - \mathbf{p}_i)}{\|\mathbf{p} - \mathbf{p}_i\|}$ is the separation distance unit vector (see also Fig. 2). Thus, the interaction force acting on the magnetic particle 2 with dipole moment ${}^2 \mathbf{m}$ due to the presence of a particle 1 with dipole moment ${}^1 \mathbf{m}$ is expressed as [24]:

$$\begin{aligned} \mathbf{f}_{\text{im},2 \leftarrow 1} = & \frac{3\mu_0}{4\pi d_{12}^4} \left(\mathbf{d} ({}^1 \mathbf{m} \cdot {}^2 \mathbf{m}) + {}^1 \mathbf{m} (\mathbf{d}_{12} \cdot {}^2 \mathbf{m}) \right. \\ & \left. + {}^2 \mathbf{m} (\mathbf{d}_{12} \cdot {}^1 \mathbf{m}) \right. \\ & \left. - 5 \frac{\mathbf{d}}{d_{12}^2} (\mathbf{d}_{12} \cdot {}^2 \mathbf{m}) (\mathbf{d}_{12} \cdot {}^1 \mathbf{m}) \right) \quad (7) \end{aligned}$$

where $\mathbf{d}_{12} = d_{12} \mathbf{d}_0$ is the separation distance vector between the two particles, as illustrated in Fig. 2. Obviously, the Newton's third law implies: $\mathbf{f}_{\text{im},1 \leftarrow 2} = -\mathbf{f}_{\text{im},2 \leftarrow 1}$. Then, the magnetic interaction force magnitude between the particles 1 and 2 is defined as :

$$\begin{aligned} {}^{12} \mathbf{f}_{\text{im}} = & \|\mathbf{f}_{\text{im},1 \leftarrow 2}\| = \|\mathbf{f}_{\text{im},2 \leftarrow 1}\| \\ = & \frac{3\mu_0 \|\mathbf{m}_1 \mathbf{m}_2\|}{4\pi d^4} \sqrt{1 - 2 \cos^2 \theta_{12} + 5 \cos^4 \theta_{12}} \quad (8) \end{aligned}$$

with θ_{12} the angle between the dipole moment \mathbf{m}_j ($j=1$ or 2) and the separation distance direction \mathbf{d}_{12} (see Fig. 2).

First, we assumed that all dipole moments of the magnetic boluses are aligned along the uniform field \mathbf{b}_0 and are not saturated to their maximum value $\mathbf{m}_{j,\text{sat}}$. Secondly, for a given separation distance d_{12} the above interaction force magnitude is minimal for $\theta_{\text{min}} = k180^\circ + \{63.43^\circ; 116.56^\circ\}$, and maximal for $\theta_{\text{max}} = k180^\circ + \{0; 180^\circ\}$, $\forall k \in \mathbb{Z}$. Commonly, the magnetic gradient is limited by the capability of the magnetic coils system. Hence, the minimal controllable separation distance is given by:

$$d_{\text{min}} = \left(\frac{6\mu_0 \|\mathbf{m}_a \mathbf{m}_b\|}{4\pi \|\mathbf{f}_{\text{m,max}}\|} \right)^{1/4} \quad (9)$$

with $\|\mathbf{f}_{\text{m,max}}\|$ the maximum magnetic force induced by the magnetic coils system.

Let us consider two magnetic boluses with different radius, $r_1 = 1075 \mu\text{m}$, $r_2 = 805 \mu\text{m}$. The magnetic interaction force \mathbf{f}_{im} is computed for different magnetic field \mathbf{b} strengths which is used to magnetize the boluses. As shown in the Fig. 3, the magnetic interaction forces is in the order of micro-newtons for a separation distances in the range of [3.5 mm; 7 mm]. For a separation distance $d = 3.5$ mm, the magnetic interaction

force is about $1.7 \mu\text{N}$ of magnitude when an external magnetic field of $\|\mathbf{b}_0\| = 35 \text{ mT}$ is applied.

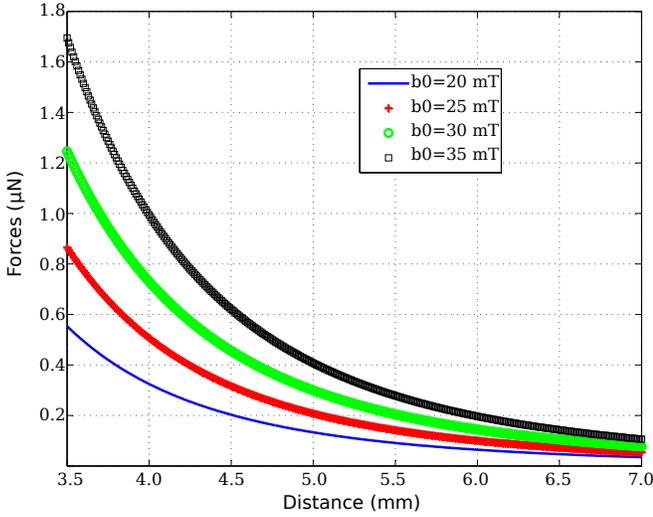


Figure 3. The magnetic interaction force of the bolus as function a separation distance.

The magnetic interaction forces have been theoretically determined. It remains to investigate the other interaction forces acting between boluses.

C. Total Interaction Forces

The considered magnetic boluses are ferrofluid droplet immersed in a microfluidic environment. As previously presented, in the presence of a magnetic field, the boluses moved through the magnetic interaction force \mathbf{f}_{im} acting between them. Other non-magnetic interaction forces also appear, and the total interaction forces could be formulated as follows:

$$\begin{aligned} {}^{12}\mathbf{f}_i &= {}^{12}\mathbf{f}_{\text{im}} + {}^{12}\mathbf{f}_{\text{in}} \\ {}^{21}\mathbf{f}_i &= {}^{21}\mathbf{f}_{\text{im}} + {}^{21}\mathbf{f}_{\text{in}} \end{aligned} \quad (10)$$

where \mathbf{f}_{in} embeds all other non-magnetic interaction force (non-contact force). Moreover, when magnetic boluses navigate in a microfluidic environment they are subject to the hydrodynamic drag force \mathbf{f}_{d} . Considering a low motion (that is at a low Reynold number), these hydrodynamic drag forces are estimated from the Stoke's law:

$${}^i\mathbf{f}_{\text{d}} = -6\pi\eta_f r_i \mathbf{v}_i \quad (11)$$

with η_f the fluid viscosity; r_i and \mathbf{v}_i respectively the radii and velocities of the i th magnetic bolus. Thus, using the second principle of Newton's law implies the following dynamic on two boluses:

$$\begin{aligned} m_1 \ddot{\mathbf{x}}_1 &= {}^1\mathbf{f}_{\text{d}} + {}^{12}\mathbf{f}_i \\ m_2 \ddot{\mathbf{x}}_2 &= {}^2\mathbf{f}_{\text{d}} + {}^{21}\mathbf{f}_i \end{aligned} \quad (12)$$

where $\ddot{\mathbf{x}}_i$ and m_i are the acceleration and the mass of the i th bolus. Therefore, from (10) and (12), the total interaction forces could be estimated from:

$$\begin{aligned} {}^{12}\mathbf{f}_i &= m_1 \ddot{\mathbf{x}}_1 - {}^1\mathbf{f}_{\text{d}} \\ {}^{21}\mathbf{f}_i &= m_2 \ddot{\mathbf{x}}_2 - {}^2\mathbf{f}_{\text{d}} \end{aligned} \quad (13)$$

III. EXPERIMENTAL RESULTS

A. Experimental setup

To estimate experimentally the magnetic interaction forces between two boluses experiments have conducted on our experimental testbed [25]. The experimental platform comprise an electromagnetic actuation (EMA) system developed by Aeon Scientific™. The EMA system consists of three nested sets of Maxwell coils and one nested set of Helmholtz coils [25], and is illustrated in Fig.4(a). Such arrangement allows generating a constant-gradient magnetic field pointing in x , y , and z -axis directions. The generated uniform magnetic field strength is limited to $\|\mathbf{b}_{\text{max}}\| = 30 \text{ mT}$. Magnetic field will thus be exerted to magnetize the different magnetic boluses that are placed inside a tank filled with an aqueous solution of 50 % glycerin that is closed to the blood viscosity Fig.4(b). The Table I summarizes the relevant experimental parameters set used for the magnetic boluses.

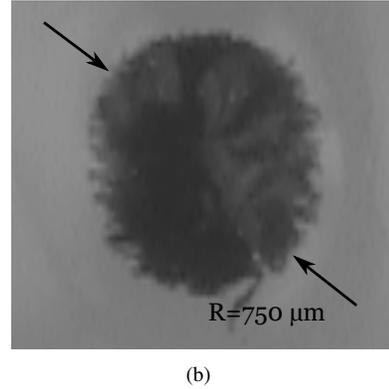
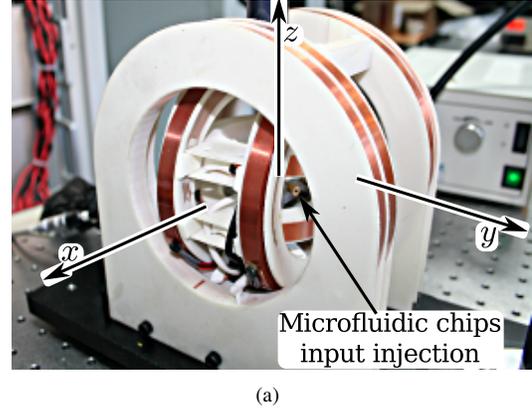


Figure 4. Experimental setup: (a) 3D Maxwell-Helmholtz coils and (b) a bolus in a tank.

B. Magnetic Bolus Preparation

The magnetic microrobot is a colloidal suspension of superparamagnetic iron oxide particles (SPIO). More precisely, the ferrofluid is a heterogeneous mixture consisting of a colloidal liquid where magnetite particles (Fe_3O_4) are suspended in a carrier fluid, that will be in the end replaced with the therapeutic agent. Therefore, two phases can be

Table I
EXPERIMENTAL PARAMETERS SET

Parameters	Values
χ	0.901
τ_m	0.6
ρ_f	3488 (kg/m ³)
$\rho_{f/b}$	920 (kg/m ³)
ρ_p	5242 (kg/m ³)
r_1	1075 (μ m)
r_2	805 (μ m)
η_f	4 (mPa.s)
$\eta_{f/b}$	62 (mPa.s)

differentiated: the carrier phase and the surrounding liquid. To mimic different types of drug load, here, SPIO powder (Inoxia Ltd, UK) are suspended in oleaginous phase, and is shown on the Fig.5(a). Sunflower oil, which is biocompatible, biodegradable, environmentally friendly, was selected as the oleaginous carrier liquid in the present work. Thus, the bolus is an homemade ferrofluid that is prepared by mixing sunflower seed oil and the SPIO powder for 30 min in a volumetric ratio of $\tau_m = 60\%$ magnetic particles. This chosen τ_m ratio is related to the optimal magnetization rate that allows conveying an optimum drug load [17]. Then the density of the magnetic boluses is basically computed from:

$$\rho_b = (\tau_m \rho_p + \rho_{f/b}(1 - \tau_m)) \quad (14)$$

with τ_m the magnetization rate; ρ_p the density of the SPIO (Fe₃O₄) particles; $\rho_{f/b}$ the ferrofluid carrier fluid density (see also Table I).

C. Magnetic interaction force estimation

We evaluated experimentally the magnetic interaction forces between two boluses along a similar direction (x-axis) in a tank filled with an aqueous solution of 50 % glycerin. As shown in Fig. 5(a), the initial distance separating the boluses is settled to $d = 6.75$ mm where the magnetic attractive forces could be neglected. Then, each magnetic bolus experiences the same magnetic field $\mathbf{b} = (30 \text{ mT}, 0, 0)^T$. The bolus B_1 and B_2 move from their initial positions $x_1 = -1.7$ mm, $x_2 = 5.05$ mm respectively to the final positions $x_{1f} = 0.78$ mm, $x_{2f} = 4.34$ mm depicted in Fig. 5(b). As expected, this one dimensional motion is caused by the presence of an interaction force acting between the two boluses.

As shown in Fig.6, the separation distance between the boluses decreases as function the time until reaching approximately 3.5 mm which corresponds to the collision between the boluses (cf. Fig.5(b)). The experimental magnetic interaction force acting between boluses is given in Fig.7. A maximal value of $1.2 \mu\text{N}$ is obtained, just before contact between the boluses. Obviously, this magnetic interaction force decreases as a separation distance increases.

The measured mean velocities of boluses are $\mathbf{v}_1 = 1.036$ mm/s and $\mathbf{v}_2 = -2.54$ mm/s; and the corresponding accelerations $\ddot{\mathbf{x}}_1 = -0.0083$ mm/s² and $\ddot{\mathbf{x}}_2 = 0.0979$ mm/s². The magnetic boluses interaction forces strength is estimated

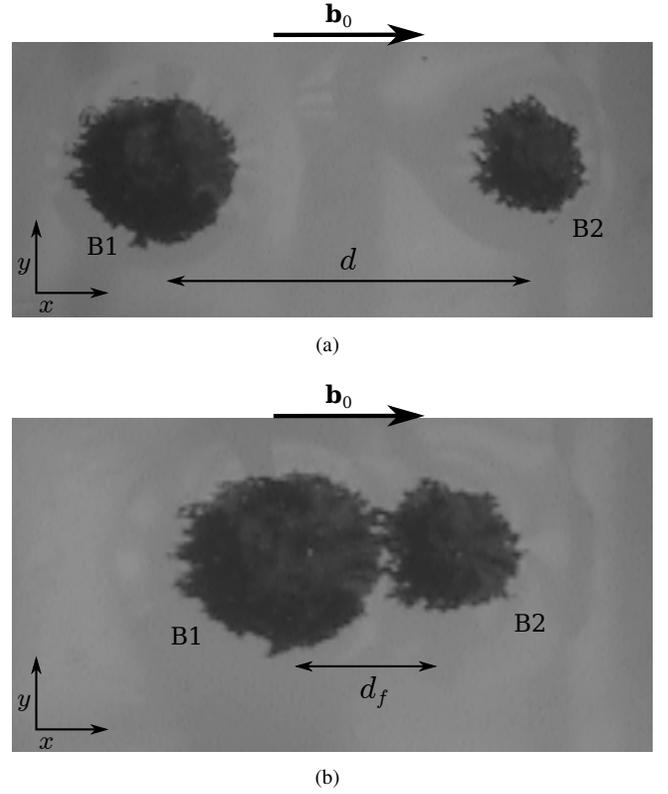


Figure 5. Two magnetic boluses immersed into a Water/Glycerol tank: (a) the boluses are in their initial positions, (b): the boluses are in collision after application of magnetic field $b_{0x} = 30$ mT.

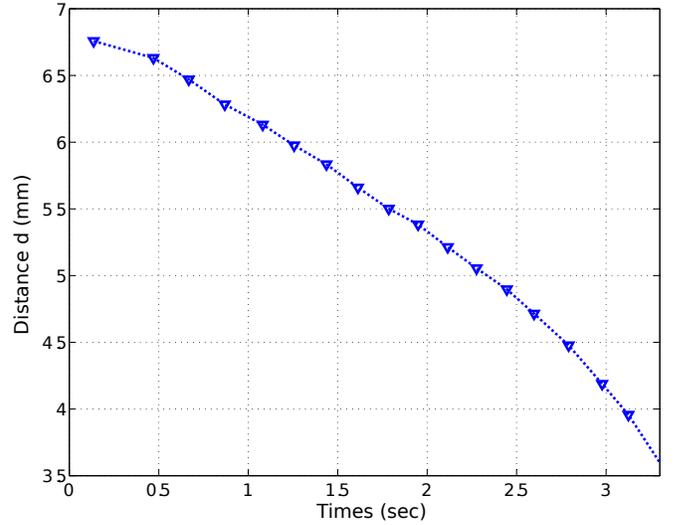


Figure 6. Evolution of separation distance d as function the time.

about ${}^{12}\mathbf{f}_i = 0.9435 \mu\text{N}$ and ${}^{21}\mathbf{f}_i = 3.0904 \mu\text{N}$. The non-magnetic interaction forces given in Fig.8 are in the order of few micro-newton. Therefore, these non-magnetic forces are not negligible [26]. Especially, these interaction forces increase with the boluses' separation distance d .

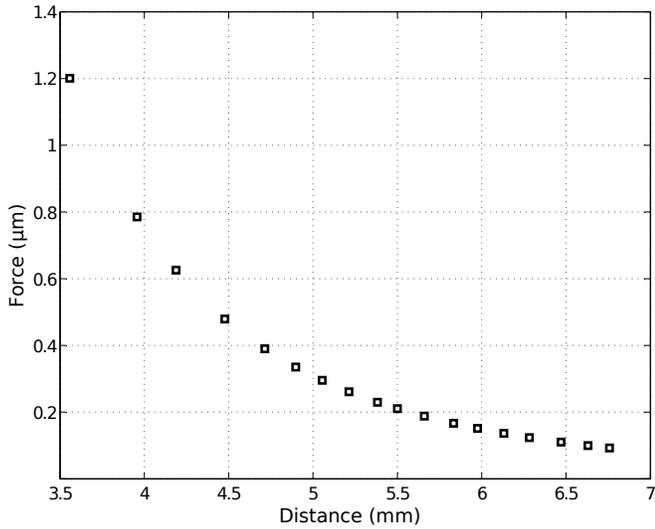


Figure 7. The experimental magnetic interaction force between boluses as function of the separation distance.

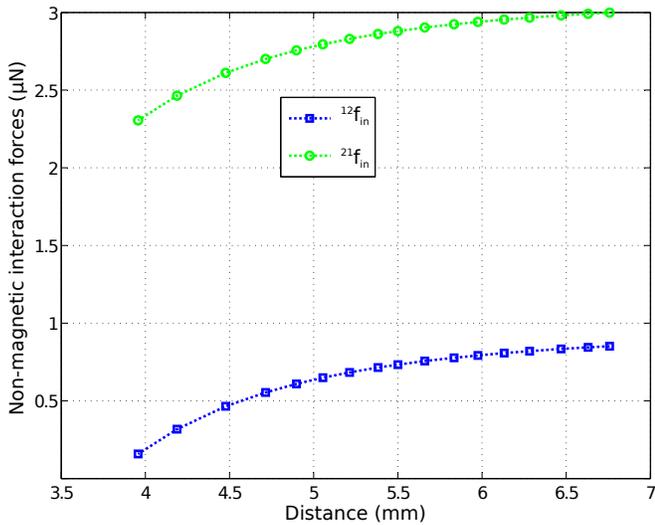


Figure 8. The experimental non magnetic interaction forces between boluses as function a separation distance.

D. Discussions

From the experimental results, it appears that the interaction force experienced on the bolus B_2 (the small one) due to the presence of the bolus B_1 (the big one) is greater than the interaction force experienced by the bolus B_1 due to the presence of the bolus B_2 . These difference are related to their shape and volume difference. Obviously, each magnetic bolus moves thanks to the magnetic interaction, but experience different drag forces f_{d1} and f_{d2} that are related to their shape. Secondly, the non-magnetic interaction forces are mainly related to the presence of hydrodynamic gradient pressure [26]. Actually, the motion of the magnetic bolus B_1 induce a high gradient pressure around the bolus B_2 which generates a non-contact forces. This gradient pressure is much higher around the bolus B_2 , than in the vicinity of the bolus B_1 .

IV. CONCLUSION

This paper described the theoretical formulation of the magnetic interaction forces between two boluses. From experiments, we have estimated the magnetic and non magnetic interaction forces. This estimation show that the boluses experienced another interaction forces than magnetic interactions. The non magnetic forces is in micro newtons order and are not negligible. These forces are caused by the presence of a high gradient pressure around the bolus. When the boluses differ in the size, the forces experienced on the one bolus due to the presence of another bolus differ. In the future work, we develop a robust control law in order to control the motion of the boluses with the presence of the interaction forces.

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